
Learning Scalar Fields for Molecular Docking with Fast Fourier Transforms

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Abstract

Molecular docking is critical to structure-based virtual screening, yet the throughput of such workflows is limited by the expensive optimization of scoring functions involved in most docking algorithms. We explore how machine learning can accelerate this process by learning a scoring function with a functional form that allows for more rapid optimization. Specifically, we define the scoring function to be the cross-correlation of multi-channel ligand and protein scalar fields parameterized by equivariant graph neural networks, enabling rapid optimization over rigid-body degrees of freedom with fast Fourier transforms. Moreover, the runtime of our approach can be amortized at several levels of abstraction, and is particularly favorable for virtual screening settings with a common binding pocket. We benchmark our scoring functions on two simplified docking-related tasks: decoy pose scoring and rigid conformer docking. Our method attains similar but faster performance on crystal structures compared to the Vina and Gnina scoring functions, and is more robust on computationally predicted structures.

1 Introduction

Proteins are the macromolecular machines that drive almost all biological processes, and much of early-stage drug discovery focuses on finding molecules which bind to and modulate their activity. *Molecular docking*—the computational task of predicting the binding pose of a small molecule to a protein target—is an important step in this pipeline. Traditionally, docking has been formulated as an optimization problem over a *scoring function* designed to be a computational proxy for the free energy (Torres et al., 2019; Fan et al., 2019). Such scoring functions are typically a sum of pairwise interaction terms between atoms with physically-inspired functional forms (Quiroga & Villarreal, 2016). While these terms are simple and hence fast to evaluate, exhaustive sampling or optimization over the space of ligand poses is difficult and leads to the significant runtime of docking software.

ML-based scoring functions for docking have been an active area of research (Yang et al., 2022; Crampon et al., 2022), largely aiming to more accurately model the free energy based on a docked pose for downstream screening. However, they have not addressed nor reduced the computational cost required to produce these poses in the first place. In this work, we explore a different paradigm and motivation for machine learning scoring functions, with the specific aim of *accelerating scoring and optimization* of ligand poses for high-throughput molecular docking. To do so, we forego the physics-inspired functional form of commonly used scoring functions, and instead frame the problem as that of learning *scalar fields* independently associated with the 3D structure of the protein and ligand, respectively. We then define the score to be the cross-correlation between the overlapping scalar fields when oriented according to the ligand pose. While seemingly more complex, these cross-correlations can be rapidly evaluated over a large number of ligand poses simultaneously using Fast Fourier Transforms (FFT) over both the translational space \mathbb{R}^3 and the rotational space $SO(3)$, allowing for significant speedups in the optimization over these degrees of freedom.

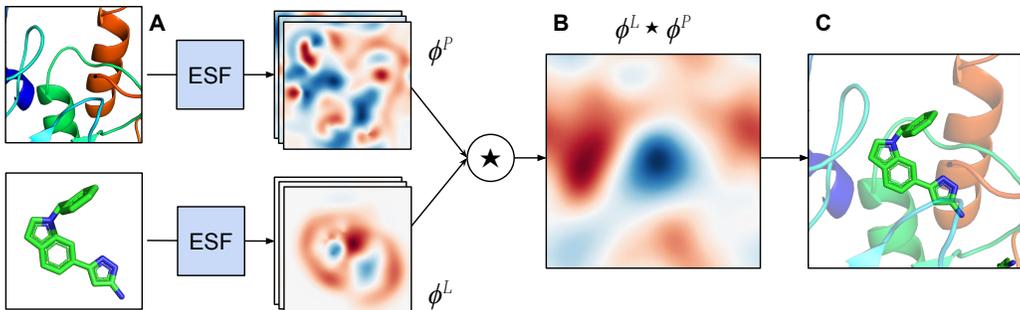


Figure 1: **Overview** of the scalar field-based scoring function and docking procedure. The translational FFT procedure is shown here; the rotational FFT is similar, albeit harder to visualize. (A) The protein pocket and ligand conformer are independently passed through equivariant scalar field networks (ESFs) to produce scalar fields. (B) The fields are cross-correlated to produce heatmaps over ligand translations. (C) The ligand coordinates are translated to the argmax of the heatmap. Additional scalar field visualizations are in Appendix E.

Further contrasting with existing ML scoring functions, the computational cost of our method can be *amortized* at several levels of abstraction, significantly accelerating runtimes for optimized workflows. For example, unlike methods that require one neural network forward pass per pose, our network is evaluated once per protein structure or ligand conformer *independently*. Post-amortization, we attain translational and rotational optimization runtimes as fast as 160 μs and 650 μs , respectively, with FFTs. Such throughputs, when combined with effective sampling and optimization, could make docking of very large compound libraries feasible with only modest resources.

2 Method

Equivariant Scalar Fields We consider the inputs to a molecular docking problem to be a pair of protein structure and ligand molecule, encoded as a featurized graphs G^P, G^L , and with the protein structure associated with alpha carbon coordinates $\mathbf{X}^P = [\mathbf{x}_1^P, \dots, \mathbf{x}_{N_P}^P] \in \mathbb{R}^{3 \times N_P}$. The molecular docking problem is to find the ligand atomic coordinates $\mathbf{X}^L = [\mathbf{x}_1^L, \dots, \mathbf{x}_{N_L}^L] \in \mathbb{R}^{3 \times N_L}$ of the true binding pose. To this end, our aim is to parameterize and learn (multi-channel) scalar fields $\phi^P := \phi(\mathbf{x}; G^P, \mathbf{X}^P)$ and $\phi^L := \phi^L(\mathbf{x}; G^L, \mathbf{X}^L)$ associated with the protein and ligand structures, respectively, such that the scoring function evaluated on any pose $\mathbf{X}^L \in \mathbb{R}^{3N_L}$ is given by

$$E(\mathbf{X}^P, \mathbf{X}^L) = \sum_c \int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}; G^P, \mathbf{X}^P) \phi_c^L(\mathbf{x}; G^L, \mathbf{X}^L) d^3\mathbf{x} \quad (1)$$

where ϕ_c refers to the c^{th} channel of the scalar field. We introduce *equivariant scalar fields* in which the scalar field is parameterized as a sum of contributions from each ligand atom or protein alpha-carbon, each defined by its coefficients in a *spherical harmonic expansion* centered at that coordinate in 3D space. To do so, we choose a set $R_j : \mathbb{R}^+ \rightarrow \mathbb{R}$ of radial basis functions (e.g., Gaussian RBFs) in 1D and let Y_m^ℓ be the real spherical harmonics. Then we define

$$\phi_c(\mathbf{x}; G, \mathbf{X}) = \sum_{n,j,\ell,m} A_{cnj\ell m}(G, \mathbf{X}) R_j(\|\mathbf{x} - \mathbf{x}_n\|) Y_m^\ell \left(\frac{\mathbf{x} - \mathbf{x}_n}{\|\mathbf{x} - \mathbf{x}_n\|} \right) \quad (2)$$

where here (and elsewhere) we drop the superscripts L, P for common definitions. Given some constraints on how the vector of coefficients $A_{cnj\ell m}$ transforms under $SE(3)$, this parameterization of the scalar field satisfies the following important properties:

Proposition 1. *Suppose the scoring function is parameterized as in Equation 2 and for any $R \in SO(3), \mathbf{t} \in \mathbb{R}^3$ we have $A_{cnj\ell m}(G, R \cdot \mathbf{X} + \mathbf{t}) = \sum_{m'} D_{mm'}^\ell(R) A_{cnj\ell m'}(G, \mathbf{X})$ where $D^\ell(R)$ are the (real) Wigner D-matrices, i.e., irreducible representations of $SO(3)$. Then for any $g \in SE(3)$,*

1. *The scalar field transforms equivariantly: $\phi_c(\mathbf{x}; G, g \cdot \mathbf{X}) = \phi_c(g^{-1} \cdot \mathbf{x}; G, \mathbf{X})$.*
2. *The scoring function is invariant: $E(g \cdot \mathbf{X}^P, g \cdot \mathbf{X}^L) = E(\mathbf{X}^P, \mathbf{X}^L)$.*

See Appendix C for the proof. We choose to parameterize $A_{cnj\ell m}(G, R, \mathbf{X})$ with E3NN graph neural networks (Thomas et al., 2018; Geiger & Smidt, 2022), which are specifically designed to satisfy these equivariance properties and produce all coefficients in a single forward pass. The core of our method consists of the training of two such equivariant scalar field networks (ESFs), one for the ligand and one for the protein, which then parameterize their respective scalar fields. Next, we show how this parameterization enables ligand poses related by rigid body motions to some reference pose to be rapidly evaluated with fast Fourier transforms (all derivations in Appendix C).

FFT Over Translations. Given some reference pose \mathbf{X}^L , the score as a function of the translation is just the cross-correlation of the fields ϕ^L and ϕ^P :

$$E(\mathbf{X}^P, \mathbf{X}^L + \mathbf{t}) = \sum_c \int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}) \phi_c^L(\mathbf{x} - \mathbf{t}) d^3 \mathbf{x} = \sum_c (\phi_c^L \star \phi_c^P)(\mathbf{t}) \quad (3)$$

where we have dropped the dependence on G, \mathbf{X} for cleaner notation and applied Proposition 1. By the convolution theorem, these cross-correlations may be evaluated using Fourier transforms $\phi_c^L \star \phi_c^P = \frac{1}{(2\pi)^{3/2}} \mathcal{F}^{-1} \left\{ \mathcal{F}[\phi_c^L] \cdot \mathcal{F}[\phi_c^P] \right\}$. Hence, in order to simultaneously evaluate all possible translations of the reference pose, we need to compute the Fourier transforms of the protein and ligand scalar fields. Conveniently, the functional form of the scalar field allows us to immediately obtain the Fourier transform via the expansion coefficients $A_{cnj\ell m}$:

$$\mathcal{F}[\phi_c](\mathbf{k}) = \sum_n e^{-i\mathbf{k} \cdot \mathbf{x}_n} \sum_{\ell} (-i)^\ell \sum_{m,n} A_{cnj\ell m} \mathcal{H}_\ell[R_j](\|\mathbf{k}\|) Y_\ell^m(\mathbf{k}/\|\mathbf{k}\|) \quad (4)$$

where now Y_ℓ^m must refer to the complex spherical harmonics and the coefficients must be transformed correspondingly, and $\mathcal{H}_\ell[R_j](k) = \sqrt{\frac{2}{\pi}} \int_0^\infty j_\ell(kr) R_j(r) r^2 dr$ is the ℓ^{th} order spherical Bessel transform of the radial basis functions.

FFT Over Rotations. Suppose that given some reference pose \mathbf{X}^L , the protein and ligand scalar fields are both expanded around some common coordinate system origin using the complex spherical harmonics and a set of *global radial basis functions* $S_j(r)$:

$$\phi_c(\mathbf{x}) = \sum_{j,\ell,m} B_{cj\ell m} S_j(\|\mathbf{x}\|) Y_\ell^m(\mathbf{x}/\|\mathbf{x}\|) \quad (5)$$

We seek to simultaneously evaluate the score of poses generated via rigid rotations of the ligand, which (thanks again to Proposition 1) is given by the rotational cross-correlation

$$E(\mathbf{X}^P, R, \mathbf{X}^L) = \sum_c \int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}) \phi_c^L(R^{-1}\mathbf{x}) d^3 \mathbf{x} \quad (6)$$

Cross-correlations of this form have been previously studied for rapid alignment of crystallographic densities (Kovacs & Wriggers, 2002) and of signals on the sphere in astrophysics (Wandelt & Górski, 2001). It turns out that they can also be evaluated in terms of Fourier sums:

$$\int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}) \phi_c^L(R^{-1}\mathbf{x}) d^3 \mathbf{x} = \sum_{\ell,m,h,n} d_{mh}^\ell d_{hn}^\ell I_{mn}^\ell e^{i(m\xi+h\eta+n\omega)} \quad (7)$$

where ξ, η, ω are related to the the Euler angles of the rotation R , d^ℓ is the (constant) Wigner D -matrix for a rotation of $\pi/2$ around the y -axis, and $I_{mn}^\ell = \sum_{j,k} B_{cj\ell m}^P \overline{B_{ck\ell n}^L} G_{jk}$ where $G_{jk} = \int_0^\infty S_j(r) S_k(r) r^2 dr$. Thus the main remaining task is to compute the complex coefficients $B_{cj\ell m}$ of the ligand and protein scalar fields, respectively, given the expansions in "local" radial and spherical harmonic bases. It turns out that this can be done via a linear transformation of the local expansion coefficients, the details of which are deferred to Appendix B.

Inference. We now examine how the rapid cross-correlation procedures presented thus far are used in inference (the training procedure is deferred to Appendix B). Given a rigid protein structure \mathbf{X}^P , the high-level task is to score or optimize candidate ligand poses \mathbf{X}^L . A large variety of possible workflows can be imagined; however, for proof of concept and for our experiments in Section 3 we describe and focus on the following relatively simple inference workflows (Appendix B and D):

Table 1: **Typical runtimes** of the computations involved in inference-time scoring and optimization procedures, measured on PDBBind with one V100 GPU.

Frequency	Computation	TF	RF	TS	RS	Runtime
Per protein structure	Coefficients $A_{cnj\ell m}$	✓	✓	✓	✓	65 ms
	FFT coefficients	✓		✓		7.0 ms
	↔ Per translation		✓		✓	80 ms
Per ligand conformer	Coefficients $A_{cnj\ell m}$	✓	✓	✓	✓	4.3 ms
	Global expansion $B_{cj\ell m}$		✓		✓	17 ms
	↔ Per rotation	✓		✓		1.6 ms
Per conformer × rotation	Translational FFT	✓				160 μ s
Per conformer × translation	Rotational FFT		✓			650 μ s
Per pose	Translational scoring			✓		1.0 μ s
	Rotational scoring				✓	8.2 μ s

- Given a ligand conformer, we use either **translational FFT (TF)** or **rotational FFT (RF)** to optimize a pose over the corresponding degree of freedom, and conduct a grid-based search over the other degree of freedom. We return the best-scoring pair as the docked pose.
- Given a list of poses, we score them by computing the (translational or rotational) Fourier coefficients of the protein and ligand and summing their products in Fourier space. We call this procedure **translational scoring (TS)** and **rotational scoring (RF)**, respectively.

The runtime of these workflows can vary significantly depending on the parameters, i.e., number of proteins, ligands, conformers, rotations, and translations, with amortizations possible at several levels (Table 1). We highlight that the **RF** workflow is well-suited for virtual screening since the precomputations for the protein and ligand translations within a pocket can be amortized across all ligands. Furthermore, if the ligands are drawn from a shared library, their coefficients can also be precomputed independent of any protein, leaving only the rotational FFT as the cost per ligand-protein pair. Thus our method can lend itself to the engineering of very high-throughput workflows.

3 Experiments

We train and test our model on the PDBBind dataset (Liu et al., 2017) with splits as defined by Stärk et al. (2022). We train two variants of our model: ESF and ESF-N, where the latter is trained with rotational and translational noise injected into the examples to increase model robustness. For evaluation, we consider both the crystal structures in the PDBBind test split and their counterpart ESMFold complexes. We also collect a test set of 77 crystal structures of phosphodiesterase 10A (PDE10A) with different ligands bound to the same pocket (Tosstorff et al., 2022), which benchmarks the benefits of our runtime amortization. To compare against baselines, we note that a scoring function by itself is not directly comparable to complete docking programs, which also include tightly integrated conformer search, pose clustering, and local refinement algorithms. To focus on the development of the *scoring function* itself, we consider two simplified evaluation settings: (1) **scoring decoy poses** with the aim of identifying the best pose among them, and (2) **docking rigid conformers** to a given pocket, similar to the re-docking setup in Stärk et al. (2022). We select Gnina (McNutt et al., 2021) as the baseline docking software. For the scoring function, we evaluate Gnina’s namesake CNN (Ragoza et al., 2017) and the traditional scoring function of Vina (Trott & Olson, 2010). Both scoring functions are widely used and are natively supported by the Gnina program.

Scoring Decoys For each PDBBind test complex, we generate $32^3 - 1 = 32767$ decoy poses by sampling 31 translational, rotational, and torsional perturbations to the ground truth pose and considering all their possible combinations (details in Appendix F). The quality of each scoring function is evaluated with the AUROC when used as a $<2\text{\AA}$ RMSD classifier, the RMSD of the top-ranked pose (Top RMSD), the rank of the lowest-RMSD pose (Top Rank), and the fraction of selected poses under 2\AA RMSD. As shown in Table 2, our method is competitive with the Gnina and Vina scoring functions on crystal structures and better on ESMFold structures. In terms of runtime per pose, our method is faster than Vina by several orders of magnitude, with even greater acceleration compared to the Gnina CNN. The runtime improvement per complex is more tempered since the different proteins and ligand in every complex limit the opportunity for amortization. In fact, of the

Table 2: **Decoy scoring results.**

Method	Crystal structures				ESMFold structures				Time per	
	<2 Å AUROC	Top RMSD	Top Rank	% <2 Å	<2 Å AUROC	Top RMSD	Top Rank	% <2 Å	Pose	Complex
Vina	0.93	0.54	2	91	0.86	2.43	419	43	3.4 ms	110 s
Gnina	0.90	0.59	3	83	0.84	2.19	1110	46	13.0 ms	426 s
ESF-TS	0.87	<u>0.59</u>	<u>3</u>	<u>87</u>	0.82	1.38	24	57	1.0 μ s	3.2 s
ESF-RS	0.87	0.63	<u>3</u>	85	0.82	1.75	22	53	8.2 μ s	5.7 s
ESF-N-TS	<u>0.92</u>	0.69	4	81	0.87	1.64	22	54	1.0 μ s	3.2 s
ESF-N-RS	<u>0.92</u>	0.75	5	80	0.87	1.74	26	53	8.2 μ s	5.7 s

Table 3: **Rigid conformer docking results.** The median RMSD of our method (ESF) is lower-bounded at 0.5–0.6 Å by the resolution of the search grid (Appendix G). The runtime is shown as an average per complex, excluding / including pre-computations that can be amortized.

Method	PDBBind test							
	Crystal		ESMFold		Runtime	PDE10A		
	% <2 Å	Med. RMSD	% <2 Å	Med. RMSD		% <2 Å	Med. RMSD	Runtime
Vina	79	0.32	24	6.1	20 s	74	0.75	6.1 s
Gnina	77	0.33	28	5.9	23 s	73	0.77	6.0 s
ESF-TF	70	1.13	31	4.6	0.8 s / 8.3 s	67	1.20	1.0 s / 7.1 s
ESF-RF	71	<u>0.97</u>	32	4.4	0.5 s / 67 s	<u>73</u>	<u>0.82</u>	0.5 s / 1.5 s
ESF-N-TF	72	1.10	46	2.9	0.7 s / 8.2 s	64	1.11	1.0 s / 7.2 s
ESF-N-RF	<u>73</u>	1.00	47	3.0	0.5 s / 68 s	70	1.00	0.5 s / 1.5 s

total runtime per complex, only 1% (**TS**) to 5% (**RS**) is due to the pose scoring itself, with the rest due to preprocessing that must be done for every new protein and ligand independently. Hence, the total possible runtime improvement per complex is significantly greater for more suitable workflows.

Docking Conformers We consider the task of pocket-level docking where all methods are given as input the ground-truth conformer and a binding pocket defined with 8 Å of translational uncertainty. As shown in Table 3, both our method and the baselines obtain high success rates on PDBBind crystal structures. On ESMFold structures, however, our method obtains nearly twice the success rate (47% vs 28%) of the baseline scoring functions. Because of the nature of the PDBBind workflow, the total runtime is comparable to or slower than the baselines when precomputations are taken into account. However, in terms of the pose optimization itself, our method is significantly faster than the baselines, despite performing a brute force search over the non-FFT degrees of freedom. While it is also possible to trade-off performance and runtime by changing various Gnina settings from their default values, our method expands the Pareto front currently available with the Gnina pipeline (Appendix G; Figure 6). To more concretely demonstrate the runtime improvements of our method with amortization, we then dock the conformers in the PDE10A dataset where, because of the common pocket, all protein coefficients are computed only once. For the **RF** procedure in particular, the amortization of these coefficients leads to a 45x speedup in the overall runtime (67 s \rightarrow 1.5 s).

4 Conclusion

We have proposed a machine-learned based scoring function for accelerating pose optimization in molecular docking. Our scoring function shows comparable performance but improved runtime on two docking-related tasks relative to standard baselines, and thus holds promise when integrated with other components into a full docking pipeline. These integrations may include multi-resolution search, refinement with traditional scoring functions, and architectural adaptations for conformational (i.e., torsional) degrees of freedom—all potential directions of future work. We also hope our work serves as a bridge between graph-based molecular machine learning and the literature on cross-correlations in computational structural biology and can inspire related methods for other applications.

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A Background

Molecular docking. The two key components of a molecular docking algorithm are (1) one or more scoring functions for ligand poses, and (2) a search, sampling, or optimization procedure. There is considerable variation in the design of these components and how they interact with each other, ranging from exhaustive enumeration and filtering (Shoichet et al., 1992; Meng et al., 1992) to genetic, gradient-based, or MCMC optimization algorithms (Trott & Olson, 2010; Morris et al., 1998; McNutt et al., 2021). We refer to reviews elsewhere (Ferreira et al., 2015; Torres et al., 2019; Fan et al., 2019) for comprehensive details. These algorithms have undergone decades of development and have been given rise to well-established software packages in academia and industry, such as AutoDock (Morris & Lim-Wilby, 2008), Vina (Trott & Olson, 2010) and Glide (Halgren et al., 2004). In many of these, the scoring function is designed not only to identify the binding pose, but also to predict the binding affinity or activity of the ligand (Su et al., 2018). In this work, however, we focus on learning and evaluating scoring functions for the rapid prediction of binding poses.

ML methods in docking. For over a decade, ML methods have been extensively explored to improve scoring functions for already-docked ligand poses, i.e., for prediction of activity and affinity in structural-based virtual screens (Li et al., 2021; Yang et al., 2022; Crampon et al., 2022). On the other hand, developing ML scoring functions as the direct optimization objective has required more care due the enormous number of function evaluations involved. MedusaNet (Jiang et al., 2020) and Ggina (Ragoza et al., 2017; McNutt et al., 2021) proposed to sparsely use CNNs for guidance and re-ranking (respectively) in combination with a traditional scoring function. DeepDock (Méndez-Lucio et al., 2021) used a hypernetwork to predict complex-specific parameters of a simple statistical potential. Recently, geometric deep learning models have explored entirely different paradigms for docking via direct prediction of the binding pose (Stärk et al., 2022; Zhang et al., 2022; Lu et al., 2022) or via a generative model over ligand poses (Corso et al., 2023).

FFT methods in docking. Methods based on fast Fourier transforms have been widely applied for the related problem of *protein-protein docking*. Katchalski-Katzir et al. (1992) first proposed using FFTs over the translational space \mathbb{R}^3 to rapidly evaluate poses using scalar fields that encode the shape complementarity of the two proteins. Later works extended this method to rotational degrees of freedom (Ritchie & Kemp, 2000; Ritchie et al., 2008; Padhorny et al., 2016) and additional scoring terms, such as pairwise electrostatic potentials and solvent accessibility (Gabb et al., 1997; Mandell et al., 2001; Chen & Weng, 2002). Today, FFT methods are a routine step in protein-protein docking programs such as PIPER (Kozakov et al., 2006), ClusPro (Kozakov et al., 2017), and HDock (Yan et al., 2020), where they enable the evaluation of billions of poses, typically as an initial screening step before further evaluation and refinement with a more accurate scoring function.

In contrast, FFT methods have been significantly less studied for protein-ligand docking. While a few works have explored this direction (Padhorny et al., 2018; Ding et al., 2020; Nguyen et al., 2018), these algorithms have not been widely adopted nor been incorporated into established docking software. A key limitation is that protein-ligand scoring functions are typically more complicated than protein-protein scoring functions and cannot be easily expressed as a cross-correlation between scalar fields (Ding et al., 2020). To our knowledge, no prior works have explored the possibility of overcoming this limitation by *learning* cross-correlation based scoring functions.

B Additional Methods

Global Coefficients. Here we consider the task of computing the complex coefficients $B_{cjl\ell m}$ of the ligand and protein scalar fields, respectively. This is not immediate as the fields are defined using expansions in “local” radial and spherical harmonic bases, i.e., with respect to the individual atom positions as opposed to the coordinate system origin. Furthermore, since we cannot (in practice) use a complete set of radial or angular basis functions, it is generally not possible to express the ligand or protein scalar field as defined in Equation 2 using the form in Equation 5. Instead, we propose to find the coefficients $B_{cjl\ell m}$ that give the best approximation to the true scalar fields, in the sense of least squared error.

Specifically, suppose that $\mathbf{R} \in \mathbb{R}^{N_{\text{grid}} \times N_{\text{local}}}$ are the values of N_{local} real local basis functions (i.e., different origins, RBFs, and spherical harmonics) evaluated at N_{grid} grid points and $\mathbf{A} \in \mathbb{R}^{N_{\text{local}}}$ is the vector of coefficients defining the scalar field ϕ_c . Similarly define $\mathbf{S} \in \mathbb{R}^{N_{\text{grid}} \times N_{\text{global}}}$ using the real versions of the global basis functions. We seek to find the least-squares solution $\mathbf{B} \in \mathbb{R}^{N_{\text{global}}}$ to the overdetermined system of equations $\mathbf{R}\mathbf{A} = \mathbf{S}\mathbf{B}$, which is given by

$$\mathbf{B} = (\mathbf{S}^T \mathbf{S})^{-1} \mathbf{S}^T \mathbf{R} \mathbf{A} \quad (8)$$

Notably, this is simply a linear transformation of the local coefficients $A_{cnj\ell m}$. Thus, if we can precompute the inverse Gram matrix of the global bases $(\mathbf{S}^T \mathbf{S})^{-1}$ and the inner product of the global and local bases $\mathbf{S}^T \mathbf{R}$, then for any new scalar field ϕ_c the real global coefficients are immediately available via a linear transformation. The desired complex coefficients can then be easily obtained via a change of bases. At first glance, this still appears challenging due to the continuous space of possible atomic or alpha-carbon positions, but an appropriate discretization makes the precomputation relatively inexpensive without a significant loss of fidelity.

Training. For a given training example with protein structure \mathbf{X}^P , the scoring function $E(\mathbf{X}^P, \mathbf{X}^L)$ should ideally attain a maximum at the true ligand pose $\mathbf{X}^L = \mathbf{X}^{L*}$. We equate this task to that of learning an *energy based model* to maximize the log-likelihood of the true pose under the model likelihood $p(\mathbf{X}^L) \propto \exp [E(\mathbf{X}^P, \mathbf{X}^L)]$. However, as is typically the case for energy-based models, directly optimizing this objective is difficult due to the intractable partition function.

Instead, following Corso et al. (2023), we conceptually decompose the ligand pose \mathbf{X}^L to be a tuple $\mathbf{X}^L = (\mathbf{X}^C, R, \mathbf{t})$ consisting of a zero-mean conformer \mathbf{X}^C , a rotation R , and a translation \mathbf{t} , from which the pose coordinates are obtained: $\mathbf{X}^L = R \cdot \mathbf{X}^C + \mathbf{t}$. Then consider the following *conditional* log-likelihoods:

$$\log p(\mathbf{t} \mid \mathbf{X}^C, R) = E(\mathbf{X}^P, \mathbf{X}^L) - \log \int_{\mathbb{R}^3} \exp [E(\mathbf{X}^P, R \cdot \mathbf{X}^C + \mathbf{t}')] d^3 \mathbf{t}' \quad (9a)$$

$$\log p(R \mid \mathbf{X}^C, \mathbf{t}) = E(\mathbf{X}^P, \mathbf{X}^L) - \log \int_{SO(3)} \exp [E(\mathbf{X}^P - \mathbf{t}, R' \cdot \mathbf{X}^C)] dR' \quad (9b)$$

We observe that these integrands are precisely the cross-correlations in Equations 3 and 6, respectively, and can be quickly evaluated and summed for all values of \mathbf{t}' and R' using fast Fourier transforms. Thus, the integrals—which are the marginal likelihoods $p(\mathbf{X}^C, R)$ and $p(\mathbf{X}^C, \mathbf{t})$ —are tractable and the conditional log-likelihoods can be directly optimized in order to train the neural network. Although neither technically corresponds to the joint log-likelihood of the pose, we find that these training objectives work well in practice and optimize their sum in our training procedure.

Inference. Given a rigid protein structure \mathbf{X}^P , the high-level task is to score or optimize candidate ligand poses \mathbf{X}^L . A large variety of possible workflows can be imagined; however, for proof of concept and for our experiments in Section 3 we describe and focus on the following relatively simple inference workflows (presented in greater detail in Appendix D):

- **Translational FFT (TF).** Given a conformer \mathbf{X}^C , we conduct a grid-based search over R and use FFT to optimize \mathbf{t} in order to find the best pose $(\mathbf{X}^C, R, \mathbf{t})$. To do so, we compute the Fourier coefficients (Equation 4) of the protein \mathbf{X}^P *once* and for *each* possible ligand orientation $R \cdot \mathbf{X}^C$. We then use translational cross-correlations (Equation 3) to find the best translation \mathbf{t} for each R and return the highest scoring combination.
- **Rotational FFT (RF).** Given a conformer \mathbf{X}^C , we conduct a grid-based search over \mathbf{t} and use FFT to optimize R . To do so, we compute the global expansion coefficients $B_{cjl\ell m}^P$ of

the protein $\mathbf{X}^L - \mathbf{t}$ relative to *each* possible ligand position \mathbf{t} and *once* for the ligand \mathbf{X}^C relative to its (zero) center of mass. We then use rotational cross-correlations (Equation 6) to find the best orientation R for each \mathbf{t} and return the highest scoring combination.

- **Translational scoring (TS).** Here we instead are given a list of poses $(\mathbf{X}^C, R, \mathbf{t})$ and wish to score them. Because the values of R nor \mathbf{t} may not satisfy a grid structure, we cannot use the FFT methods. Nevertheless, we can compute the (translational) Fourier coefficients of the protein \mathbf{X}^P and for each unique oriented conformer $R.\mathbf{X}^C$ of the ligand using Equation 4. We then evaluate

$$E(\mathbf{X}^P, R.\mathbf{X}^C + \mathbf{t}) = \sum_c \int_{\mathbb{R}^3} \overline{\mathcal{F}[\phi_c^P](\mathbf{k})} \cdot \mathcal{F}[\phi_c^L(\cdot; R.\mathbf{X}^C)](\mathbf{k}) \cdot e^{-i\mathbf{k}\cdot\mathbf{t}} d^3\mathbf{k} \quad (10)$$

Since the Fourier transform is an orthogonal operator on functional space, this is equal to the real-space cross-correlation.

- **Rotational scoring (RS).** Analogously, we can score a list of poses $(\mathbf{X}^C, R, \mathbf{t})$ using the global spherical expansions $B_{c_j\ell m}$. We obtain the real expansion coefficients of the protein relative to each \mathbf{t} and for each ligand conformer \mathbf{X}^C . The score for $(\mathbf{X}^C, R, \mathbf{t})$ is then given by the rotational cross-correlation

$$E(\mathbf{X}^P, R.\mathbf{X}^C + \mathbf{t}) = \sum_{c,j,k,\ell,m,n} B_{c_j\ell m}^P(\mathbf{X}^P - \mathbf{t}) B_{c_k\ell n}^L(\mathbf{X}^C) D_{mn}^\ell(R) G_{jk} \quad (11)$$

where G_{jk} is as defined previously and D_{mn}^ℓ are the real Wigner D -matrices.

C Mathematical Details

C.1 Proof of Proposition 1

Proposition 1. *Suppose the scoring function is parameterized as in Equation 2 and for any $R \in SO(3)$, $\mathbf{t} \in \mathbb{R}^3$ we have $A_{cnj\ell m}(G, R \cdot \mathbf{X} + \mathbf{t}) = \sum_{m'} D_{mm'}^\ell(R) A_{cnj\ell m'}(G, \mathbf{X})$ where $D^\ell(R)$ are the (real) Wigner D -matrices, i.e., irreducible representations of $SO(3)$. Then for any $g \in SE(3)$,*

1. *The scalar field transforms equivariantly: $\phi_c(\mathbf{x}; G, g \cdot \mathbf{X}) = \phi_c(g^{-1} \cdot \mathbf{x}; G, \mathbf{X})$.*
2. *The scoring function is invariant: $E(g \cdot \mathbf{X}^P, g \cdot \mathbf{X}^L) = E(\mathbf{X}^P, \mathbf{X}^L)$.*

Proof. Let the action of $g = (R, \mathbf{t}) \in SE(3)$ be written as $g : \mathbf{x} \mapsto R\mathbf{x} + \mathbf{t}$ and hence $g^{-1} : \mathbf{x} \mapsto R^T(\mathbf{x} - \mathbf{t})$. We first note that $\|\mathbf{x} - g \cdot \mathbf{x}_n\| = \|g^{-1} \cdot \mathbf{x} - \mathbf{x}_n\|$ and $R^T(\mathbf{x} - g \cdot \mathbf{x}_n) = g^{-1} \cdot \mathbf{x} - \mathbf{x}_n$. Then

$$\begin{aligned}
\phi_c(\mathbf{x}; G, g \cdot \mathbf{X}) &= \sum_{n,j,\ell,m} A_{cnj\ell m}(G, R \cdot \mathbf{X} + \mathbf{t}) R_j(\|\mathbf{x} - g \cdot \mathbf{x}_n\|) Y_\ell^m \left(\frac{\mathbf{x} - g \cdot \mathbf{x}_n}{\|\mathbf{x} - g \cdot \mathbf{x}_n\|} \right) \\
&= \sum_{n,j,\ell,m'} A_{cnj\ell m'}(G, \mathbf{X}) R_j(\|\mathbf{x} - g \cdot \mathbf{x}_n\|) \sum_m D_{mm'}^\ell(R) Y_\ell^m \left(\frac{\mathbf{x} - g \cdot \mathbf{x}_n}{\|\mathbf{x} - g \cdot \mathbf{x}_n\|} \right) \\
&= \sum_{n,j,\ell,m'} A_{cnj\ell m'}(G, \mathbf{X}) R_j(\|\mathbf{x} - g \cdot \mathbf{x}_n\|) Y_\ell^{m'} \left(\frac{R^T(\mathbf{x} - g \cdot \mathbf{x}_n)}{\|\mathbf{x} - g \cdot \mathbf{x}_n\|} \right) \\
&= \sum_{n,j,\ell,m'} A_{cnj\ell m'}(G, \mathbf{X}) R_j(\|g^{-1} \cdot \mathbf{x} - \mathbf{x}_n\|) Y_\ell^{m'} \left(\frac{g^{-1} \cdot \mathbf{x} - \mathbf{x}_n}{\|g^{-1} \cdot \mathbf{x} - \mathbf{x}_n\|} \right) \\
&= \phi_c(g^{-1} \cdot \mathbf{x}; G, \mathbf{X})
\end{aligned}$$

Next,

$$\begin{aligned}
E(g \cdot \mathbf{X}^P, g \cdot \mathbf{X}^L) &= \sum_c \int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}; G^P, g \cdot \mathbf{X}^P) \phi_c^L(\mathbf{x}; G^L, g \cdot \mathbf{X}^L) d^3 \mathbf{x} \\
&= \sum_c \int_{\mathbb{R}^3} \phi_c^P(g^{-1} \cdot \mathbf{x}; G^P, \mathbf{X}^P) \phi_c^L(g^{-1} \cdot \mathbf{x}; G^L, \mathbf{X}^L) d^3 \mathbf{x} \\
&= \sum_c \int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}'; G^P, \mathbf{X}^P) \phi_c^L(\mathbf{x}'; G^L, \mathbf{X}^L) d^3 \mathbf{x}'
\end{aligned}$$

where the last line has substitution $\mathbf{x}' = g^{-1} \cdot \mathbf{x}$ with g volume preserving on \mathbb{R}^3 . \square

C.2 Derivations

In this section we describe the derivations for the various equations presented in the main text. We use the following convention for the (one-dimensional) Fourier transform and its inverse:

$$\mathcal{F}[f](k) = \frac{1}{\sqrt{2\pi}} \int e^{-ikx} f(x) dx \quad (12a)$$

$$\mathcal{F}^{-1}[f](x) = \frac{1}{\sqrt{2\pi}} \int e^{ikx} f(k) dk \quad (12b)$$

Equation 4 It is well known (Wikipedia, 2023) that given a function over \mathbb{R}^3 with complex spherical harmonic expansion

$$f(\mathbf{r}) = \sum_{\ell,m} f_{\ell,m}(\|\mathbf{r}\|) Y_\ell^m(\mathbf{r}/\|\mathbf{r}\|) \quad (13)$$

its Fourier transform is given by

$$f(\mathbf{k}) = \sum_{\ell,m} (-i)^\ell F_{\ell,m}(\|\mathbf{k}\|) Y_\ell^m(\mathbf{k}/\|\mathbf{k}\|) \quad (14)$$

where

$$F_{\ell,m}(k) = \frac{1}{\sqrt{k}} \int_0^\infty \sqrt{r} f_{\ell,m}(r) J_{\ell+1/2}(kr) r dr \quad (15)$$

with J_ℓ the ℓ^{th} -order Bessel function of the first kind. Relating these to the spherical Bessel functions j_ℓ via $J_{\ell+1/2}(x) = \sqrt{2x/\pi} j_\ell(x)$, we obtain

$$F_{\ell,m}(k) = \sqrt{\frac{2}{\pi}} \int_0^\infty f_{\ell,m}(r) j_\ell(kr) r^2 dr \quad := \mathcal{H}_\ell[f_{\ell,m}](k) \quad (16)$$

To apply this to our scalar fields, we define the *translation operator* $\mathcal{T}_\mathbf{r}[f](\mathbf{x}) = f(\mathbf{x} - \mathbf{r})$ and note its composition with the Fourier transform

$$(\mathcal{F} \circ \mathcal{T}_\mathbf{r})[f] = e^{-i\mathbf{k} \cdot \mathbf{r}} \mathcal{F}[f] \quad (17)$$

We then decompose the form of our scalar fields (Equation 2) into contributions from zero-origin spherical harmonic expansions

$$\phi_c(\mathbf{x}) = \sum_n \mathcal{T}_{\mathbf{x}_n}[\phi_{cn}](\mathbf{x}) \quad (18a)$$

$$\phi_{cn}(\mathbf{x}) = \sum_{\ell,m} \underbrace{\sum_j A_{cnj\ell m} R_j(\|\mathbf{x}\|) Y_\ell^m(\mathbf{x}/\|\mathbf{x}\|)}_{\phi_{cn\ell m}(\|\mathbf{x}\|)} \quad (18b)$$

Hence, the Fourier transform of each contribution is

$$\mathcal{F}[\phi_{cn}](\mathbf{k}/\|\mathbf{k}\|) = \sum_{\ell,m} (-i)^\ell \mathcal{H}_\ell[\phi_{cn\ell m}](\|\mathbf{k}\|) Y_\ell^m(\mathbf{k}/\|\mathbf{k}\|) \quad (19)$$

Equation 4 is then obtained via Equation 17 and the linearity of the Fourier and spherical Bessel transforms.

Equation 7 We source (with some modifications) the derivation from Kovacs & Wriggers (2002). We consider the cross-correlation

$$c(R) = \int_{\mathbb{R}^3} \phi(\mathbf{x}) \overline{\psi(R^{-1}\mathbf{x})} d^3\mathbf{x} \quad (20)$$

which is the same as Equation 6 with $\phi = \phi_c^P$ and $\psi = \phi_c^L$ since ϕ_c^L is a real field. Expanding in complex spherical harmonics Y_ℓ^m and radial bases S_j :

$$\phi(\mathbf{x}) = \sum_{j,\ell,m} \Phi_{j\ell m} S_j(\|\mathbf{x}\|) Y_\ell^m(\mathbf{x}/\|\mathbf{x}\|) \quad \psi(\mathbf{x}) = \sum_{j,\ell,m} \Psi_{j\ell m} S_j(\|\mathbf{x}\|) Y_\ell^m(\mathbf{x}/\|\mathbf{x}\|) \quad (21)$$

We then obtain

$$c(R) = \sum_{j,j',\ell,\ell',m,n,m'} \overline{D_{nm'}^\ell(R)} \Phi_{j\ell m} \overline{\Psi_{j'\ell'm'}} \int_{\mathbb{R}^3} [S_j \cdot S_{j'}](\|\mathbf{x}\|) [Y_\ell^m \cdot \overline{Y_{\ell'}^{n'}}](\mathbf{x}/\|\mathbf{x}\|) d^3\mathbf{x} \quad (22a)$$

$$= \sum_{j,j',\ell,\ell',m,n,m'} \overline{D_{nm'}^\ell(R)} \Phi_{j\ell m} \overline{\Psi_{j'\ell'm'}} \underbrace{\int_0^\infty [S_j \cdot S_{j'}](r) r^2 dr}_{G_{jj'}} \underbrace{\int_{S^2} [Y_\ell^m \cdot \overline{Y_{\ell'}^{n'}}](\hat{\mathbf{r}}) d\hat{\mathbf{r}}}_{\delta_{\ell\ell'} \delta_{mn}} \quad (22b)$$

$$= \sum_{\ell,m,m'} \overline{D_{mm'}^\ell(R)} \underbrace{\sum_{j,j'} \Phi_{j\ell m} \overline{\Psi_{j'\ell'm'}} G_{jj'}}_{I_{mm'}^\ell} \quad (22c)$$

Now to evaluate the complex Wigner D -matrix, we adopt the extrinsic zyz convention for Euler angles (applied right-to-left) and note that any rotation (ϕ, θ, ψ) can be decomposed as

$$R(\phi, \theta, \psi) = R_z(\underbrace{\phi - \pi/2}_\xi) R_y(\pi/2) R_z(\underbrace{\pi - \theta}_\eta) R_y(\pi/2) R_z(\underbrace{\psi - \pi/2}_\omega) \quad (23)$$

Next, one can easily check (using the standard spherical harmonics) that the Wigner D -matrix for a rotation about the z -axis is diagonal and given by $D_{mn}^\ell(R_z(\omega)) = \delta_{mn} e^{-in\omega}$. Hence,

$$D_{mn}^\ell(R(\phi, \theta, \psi)) = e^{-im\xi} d_{mh}^\ell e^{-h\eta} d_{hn}^\ell e^{-i\omega n} \quad (24)$$

where $d^\ell = D^\ell(R_y(\pi/2))$ are constant and real. Complex conjugation then gives Equation 7.

Equation 9 The conditional likelihood is

$$\log p(\mathbf{t} \mid \mathbf{X}^C, R) = \log \frac{p(\mathbf{X}^C, R, \mathbf{t})}{p(\mathbf{X}^C, R)} \quad (25a)$$

$$= \log p(\mathbf{X}^C, R, \mathbf{t}) - \log \int_{\mathbb{R}^3} p(\mathbf{X}^C, R, \mathbf{t}') d^3 \mathbf{t}' \quad (25b)$$

$$= \log E(\mathbf{X}^P, \mathbf{X}^L) - \log \int_{\mathbb{R}^3} \exp [E(\mathbf{X}^P, R.\mathbf{X}^C + \mathbf{t}')] d^3 \mathbf{t}' \quad (25c)$$

Similarly,

$$\log p(R \mid \mathbf{X}^C, \mathbf{t}) = \log \frac{p(\mathbf{X}^C, R, \mathbf{t})}{p(\mathbf{X}^C, \mathbf{t})} \quad (26a)$$

$$= \log p(\mathbf{X}^C, R, \mathbf{t}) - \log \int_{SO(3)} p(\mathbf{X}^C, R', \mathbf{t}) dR' \quad (26b)$$

$$= \log E(\mathbf{X}^P, \mathbf{X}^L) - \log \int_{SO(3)} \exp [E(\mathbf{X}^P, R.\mathbf{X}^C + \mathbf{t})] dR' \quad (26c)$$

Finally, we move \mathbf{t} to the protein coordinates (invoking the invariance of the score E) to obtain a form consistent with the rotational cross-correlations (Equation 6).

Equation 10 Given a pose $\mathbf{X}^L = R.\mathbf{X}^C + \mathbf{t}$, we evaluate

$$E(\mathbf{X}^P, R.\mathbf{X}^C + \mathbf{t}) = \sum_c \int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}) \phi_c^L(\mathbf{x}; R.\mathbf{X}^C + \mathbf{t}) d^3 \mathbf{x} \quad (27)$$

The functional inner product is equivalent in Fourier space:

$$E(\mathbf{X}^P, R.\mathbf{X}^C + \mathbf{t}) = \sum_c \int_{\mathbb{R}^3} \overline{\mathcal{F}[\phi_c^P](\mathbf{k})} \cdot \mathcal{F}[\phi_c^L(\cdot; R.\mathbf{X}^C + \mathbf{t})](\mathbf{k}) d^3 \mathbf{k} \quad (28)$$

Then with the translation operator \mathcal{T} defined previously,

$$\phi_c^L(\mathbf{x}; R.\mathbf{X}^C + \mathbf{t}) = \mathcal{T}_{\mathbf{t}}[\phi_c(\cdot; R.\mathbf{X}^C)](\mathbf{x}) \quad (29a)$$

$$\mathcal{F}[\phi_c^L(\cdot; R.\mathbf{X}^C + \mathbf{t})](\mathbf{k}) = e^{-i\mathbf{k}\cdot\mathbf{t}} \mathcal{F}[\phi_c(\cdot; R.\mathbf{X}^C)](\mathbf{k}) \quad (29b)$$

We then substitute into Equation 28 to obtain Equation 10.

Equation 11 Given a pose $\mathbf{X}^L = R.\mathbf{X}^C + \mathbf{t}$, we assume that the field $\phi_c^P(\cdot; \mathbf{X}^P - \mathbf{t})$ and $\phi_c^L(\cdot; \mathbf{X}^C)$ are written in the real global spherical harmonic expansion:

$$\phi_c^P(\mathbf{x}; \mathbf{X}^P - \mathbf{t}) = \sum_{j,\ell,m} B_{cj\ell m}^P S_j(\|\mathbf{x}\|) Y_\ell^m(\mathbf{x}/\|\mathbf{x}\|) \quad (30a)$$

$$\phi_c^L(\mathbf{x}; \mathbf{X}^C) = \sum_{j,\ell,m} B_{cj\ell m}^L S_j(\|\mathbf{x}\|) Y_\ell^m(\mathbf{x}/\|\mathbf{x}\|) \quad (30b)$$

Then, analogously to Equation 22,

$$E(\mathbf{X}^P, R.\mathbf{X}^C + \mathbf{t}) = E(\mathbf{X}^P - \mathbf{t}, R.\mathbf{X}^C) \quad (31a)$$

$$= \sum_c \int_{\mathbb{R}^3} \phi_c^P(\mathbf{x}; \mathbf{X}^P - \mathbf{t}) \phi_c^L(R^{-1}\mathbf{x}; \mathbf{X}^C) d^3 \mathbf{x} \quad (31b)$$

$$= \sum_{c,\ell,m,m'} D_{mm'}^\ell(R) \sum_{j,j'} B_{cj\ell m}^P B_{cj'\ell m'}^L G_{jj'} \quad (31c)$$

Complex conjugation has been omitted because the coefficients and D -functions are now real.

D Algorithmic Details

Below, we present in detail the four inference procedures introduced in Section 2. The three blocks of computations are color-coded corresponding to protein preprocessing (green), ligand preprocessing (blue), and the core computation (red) and labelled with typical runtimes from Table 1 (unlabelled lines have negligible runtime). The various loop levels make clear that depending on the workflow, the protein and ligand processing precomputations can be amortized and approaches a negligible fraction of the total runtime. Note, however, that for readability we have presented the algorithms assuming that all possible combinations (i.e., of proteins, ligand conformers, rotations, and translations) are of interest; if this is not true (for example in PDBBind, or in any typical pose-scoring setting), then the full benefits of amortization may not be fully realized.

Algorithm 1: TRANSLATIONAL FFT

Input: Proteins $\{(G_i^P, \mathbf{X}_i^P)\}$, conformers $\{(G_h^L, \mathbf{X}_h^L)\}$

Output: Docked poses $(\mathbf{X}_i^P, \mathbf{X}_{ih}^L) \forall i, h$

```

foreach  $(G_i^P, \mathbf{X}_i^P)$  do // protein preprocessing
  Compute coefficients  $\mathbf{A}_i^P = \{A_{c_j n \ell m}^P(G_i^P, \mathbf{X}_i^P)\}$  with neural network; // 65 ms
  Compute Fourier-space field values  $\mathcal{F}[\phi^P]_i$  using  $\mathbf{A}_i^P, \mathbf{x}_i^P$ ; // 7.0 ms

foreach  $(G_h^L, \mathbf{X}_h^L)$  do // ligand preprocessing
  Compute coefficients  $\mathbf{A}_h^L = \{A_{c_j n \ell m}^L(G_h^L, \mathbf{X}_h^L)\}$  with neural network; // 4.3 ms
  foreach  $R_k \in \{R\}_{grid} \subset SO(3)$  do
    Compute rotated coefficients  $\mathbf{A}_{h,k}^L$  using  $D^\ell(R_k)$ ;
    Compute Fourier-space field values  $\mathcal{F}[\phi^L]_{h,k}$  using  $\mathbf{A}_{h,k}^L, R_k \mathbf{X}_h^L$ ; // 1.6 ms

foreach  $(G_i^P, \mathbf{X}_i^P)$  do // pose optimization
  foreach  $(G_h^L, \mathbf{X}_h^L)$  do
    foreach  $R_k \in \{R\}_{grid} \subset SO(3)$  do
      Compute  $E(\mathbf{X}_i^P, R_k \mathbf{X}_h^L + \mathbf{t}), \forall \mathbf{t}$  using FFT; // 160  $\mu$ s
       $E_k^*, \mathbf{t}_k^* \leftarrow \{\max, \arg \max\}_{\mathbf{t}} E(\mathbf{X}_i^P, R_k \mathbf{X}_h^L + \mathbf{t})$ ;
       $k^* \leftarrow \arg \max_k E_k^*$ ;
       $\mathbf{X}_{ih}^L \leftarrow R_{k^*} \mathbf{X}_h^L + \mathbf{t}_{k^*}^*$ ;

```

Algorithm 2: ROTATIONAL FFT

Input: Proteins $\{(G_i^P, \mathbf{X}_i^P)\}$, conformers $\{(G_h^L, \mathbf{X}_h^L)\}$

Output: Docked poses $(\mathbf{X}_i^P, \mathbf{X}_{ih}^L) \forall i, h$

```

foreach  $(G_i^P, \mathbf{X}_i^P)$  do // protein preprocessing
  Compute coefficients  $\mathbf{A}_i^P = \{A_{c_j n \ell m}^P(G_i^P, \mathbf{X}_i^P)\}$  with neural network; // 65 ms
  for  $\mathbf{t}_k \in \{\mathbf{t}\}_{grid} \subset \mathbb{R}^3$  do
    Compute global expansion  $\mathbf{B}_{i,k}^P = \{B_{c_j \ell m}\}$  from  $\mathbf{A}_i^P, \mathbf{X}_i^P - \mathbf{t}_k$ ; // 80 ms

foreach  $(G_h^L, \mathbf{X}_h^L)$  do // ligand preprocessing
  Compute coefficients  $\mathbf{A}_h^L = \{A_{c_j n \ell m}^L(G_h^L, \mathbf{X}_h^L)\}$  with neural network; // 4.3 ms
  Compute global expansion  $\mathbf{B}_h^L = \{B_{c_j \ell m}\}$  from  $\mathbf{A}_h^L, \mathbf{X}_h^L$ ; // 17 ms

foreach  $(G_i^P, \mathbf{X}_i^P)$  do // pose optimization
  foreach  $(G_h^L, \mathbf{X}_h^L)$  do
    foreach  $\mathbf{t}_k \in \{\mathbf{t}\}_{grid} \subset \mathbb{R}^3$  do
      Compute  $E(\mathbf{X}_i^P - \mathbf{t}_k, R \cdot \mathbf{X}_h^L), \forall R$  using FFT; // 650  $\mu$ s
       $E_k^*, R_k^* \leftarrow \{\max, \arg \max\}_R E(\mathbf{X}_i^P - \mathbf{t}_k, R \cdot \mathbf{X}_h^L + \mathbf{t})$ ;
       $k^* \leftarrow \arg \max_k E_k^*$ ;
       $\mathbf{X}_{ih}^L \leftarrow R_{k^*} \mathbf{X}_h^L + \mathbf{t}_{k^*}$ ;

```

Algorithm 3: TRANSLATIONAL SCORING

Input: Proteins $\{(G_i^P, \mathbf{X}_i^P)\}$, conformers $\{(G_h^L, \mathbf{X}_h^L)\}$, rotations $\{R_k\}$, translations $\{\mathbf{t}_\ell\}$ **Output:** Scores $E(\mathbf{X}_i^P, R_k \mathbf{X}_h^L + \mathbf{t}_\ell) \forall i, h, k, \ell$

```
foreach ( $G_i^P, \mathbf{X}_i^P$ ) do // protein preprocessing
  Compute coefficients  $\mathbf{A}_i^P = \{A_{c_j n \ell m}^P(G_i^P, \mathbf{X}_i^P)\}$  with neural network; // 65 ms
  Compute Fourier-space field values  $\mathcal{F}[\phi^P]_i$  using  $\mathbf{A}_i^P, \mathbf{x}_i^P$ ; // 7.0 ms
foreach ( $G_h^L, \mathbf{X}_h^L$ ) do // ligand preprocessing
  Compute coefficients  $\mathbf{A}_h^L = \{A_{c_j n \ell m}^L(G_h^L, \mathbf{X}_h^L)\}$  with neural network; // 4.3 ms
  foreach  $R_k$  do
    Compute rotated coefficients  $\mathbf{A}_{h,k}^L$  using  $D^\ell(R_k)$ ;
    Compute Fourier-space field values  $\mathcal{F}[\phi^L]_{h,k}$  using  $\mathbf{A}_{h,k}^L, R_k \mathbf{X}_h^L$ ; // 1.6 ms
foreach ( $G_i^P, \mathbf{X}_i^P$ ) do // scoring
  foreach ( $G_h^L, \mathbf{X}_h^L$ ) do
    foreach  $R_k$  do
      foreach  $\mathbf{t}_\ell$  do
        Compute  $E(\mathbf{X}_i^P, R_k \mathbf{X}_h^L + \mathbf{t}_\ell)$  using Equation 10; // 1.0  $\mu$ s
```

Algorithm 4: ROTATIONAL SCORING

Input: Proteins $\{(G_i^P, \mathbf{X}_i^P)\}$, conformers $\{(G_h^L, \mathbf{X}_h^L)\}$, rotations $\{R_k\}$, translations $\{\mathbf{t}_\ell\}$ **Output:** Scores $E(\mathbf{X}_i^P, R_k \mathbf{X}_h^L + \mathbf{t}_\ell) \forall i, h, k, \ell$

```
foreach ( $G_i^P, \mathbf{X}_i^P$ ) do // protein preprocessing
  Compute coefficients  $\mathbf{A}_i^P = \{A_{c_j n \ell m}^P(G_i^P, \mathbf{X}_i^P)\}$  with neural network; // 65 ms
  for  $\mathbf{t}_k \in \{\mathbf{t}\}_{grid} \subset \mathbb{R}^3$  do
    Compute global expansion  $\mathbf{B}_{i,k}^P = \{B_{c_j \ell m}\}$  from  $\mathbf{A}_i^P, \mathbf{X}_i^P - \mathbf{t}_k$ ; // 80 ms
foreach ( $G_h^L, \mathbf{X}_h^L$ ) do // ligand preprocessing
  Compute coefficients  $\mathbf{A}_h^L = \{A_{c_j n \ell m}^L(G_h^L, \mathbf{X}_h^L)\}$  with neural network; // 4.3 ms
  Compute global expansion  $\mathbf{B}_h^L = \{B_{c_j \ell m}\}$  from  $\mathbf{A}_h^L, \mathbf{X}_h^L$ ; // 17 ms
foreach ( $G_i^P, \mathbf{X}_i^P$ ) do // scoring
  foreach ( $G_h^L, \mathbf{X}_h^L$ ) do
    foreach  $R_k$  do
      foreach  $\mathbf{t}_\ell$  do
        Compute  $E(\mathbf{X}_i^P, R_k \mathbf{X}_h^L + \mathbf{t}_\ell)$  using Equation 11; // 8.2  $\mu$ s
```

E Learned Scalar Fields

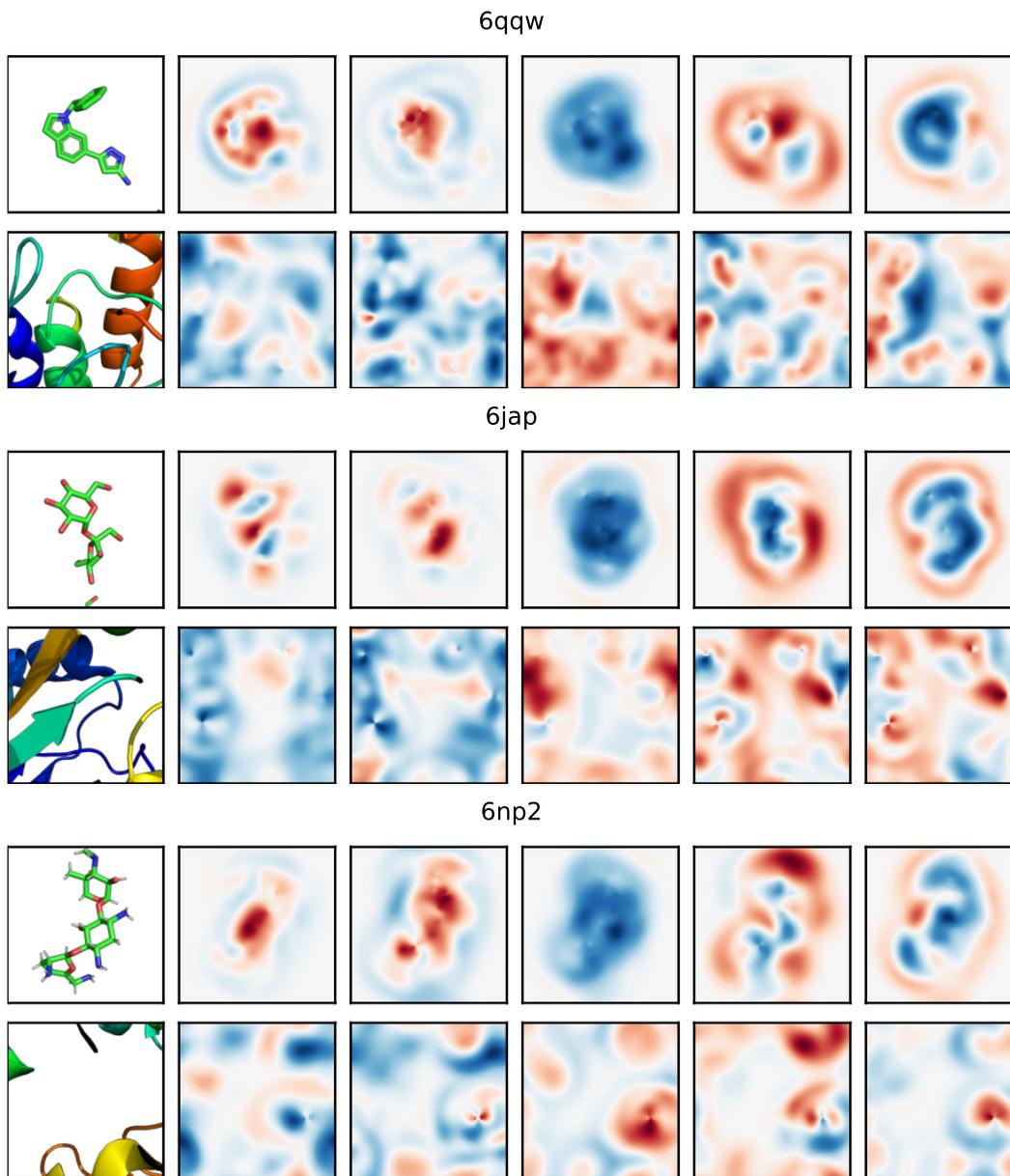


Figure 2: **Visualizations of learned scalar fields.** All five channels of the **ESF-N** learned scalar fields ϕ^L (top row) and ϕ^P (bottom row) are shown on the xy -plane passing through the center of mass of the ligand, with a box diameter of 20 Å. Positive values of the field are in **blue** and negative values in **red**. At left, the ligand and pocket structures are shown looking down the z -axis. Note that as the fields are only 2D slices, not all 3D features visible in the structures are visible in the fields.

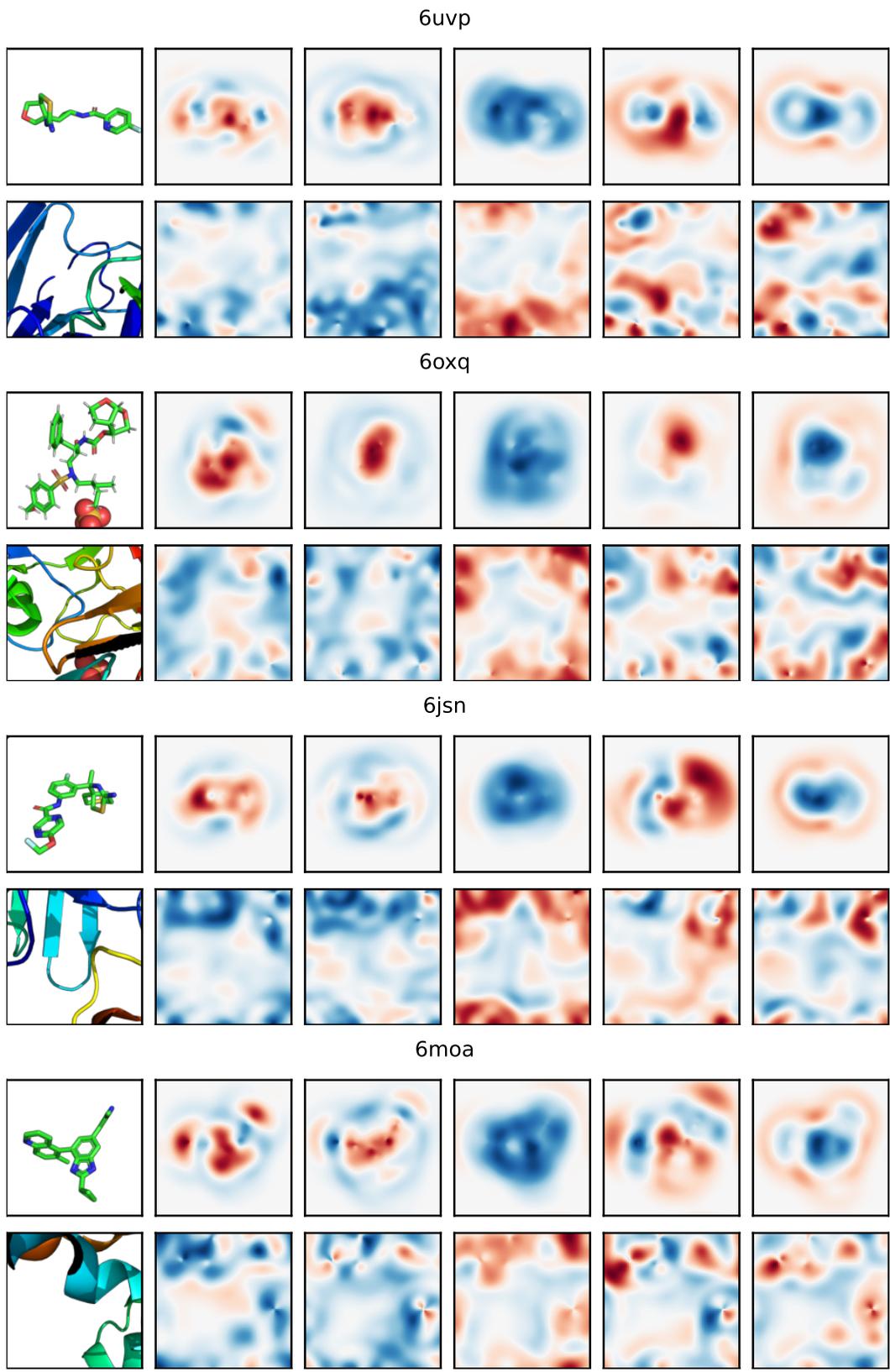


Figure 3: Visualizations of learned scalar fields, continued.

F Experimental Details

F.1 Decoy Set

Given a zero-mean ground-truth ligand pose \mathbf{X}^{L^*} , we generate $32^3 - 1 = 32767$ decoy poses via the following procedure.

- Sample 31 translational perturbations: $\mathbf{t}_i \sim \mathcal{N}(0, \mathbf{I}_3)$, $i = 1 \dots 31$ and set $\mathbf{t}_0 = \mathbf{0}$, with units in Å.
- Sample 31 rotational perturbations: $R_j = \text{FromRotvec}(\mathbf{r}_j)$, $\mathbf{r}_j \sim \mathcal{N}(0, 0.5\mathbf{I}_3)$, $j = 1 \dots 31$ and set $R_0 = \mathbf{I}_3$.
- Sample 31 noisy conformers \mathbf{X}_k^C , $k = 1 \dots 31$ by sampling torsional updates $\Delta\tau_k \sim \mathcal{N}_{\mathbb{T}}(0, (\pi/2)\mathbf{I}_m)$ where $\mathcal{N}_{\mathbb{T}}$ is a wrapped normal distribution (Jing et al., 2022) and m is the number of torsion angles. The torsional updates are applied to the smaller side of the molecule. Set $\mathbf{X}_0^C = \mathbf{X}^{L^*}$.
- Set $\mathbf{X}_{ijk}^L = R_j \mathbf{X}_k^C + \mathbf{t}_i$, $i, j, k = 0 \dots 31$ and discard $\mathbf{X}_{000}^L = \mathbf{X}^{L^*}$.

PDB ID 6A73 is excluded from the procedure due to the high level of graph symmetry and significant runtime for computing RMSDs for all decoys. Summary statistics for the decoy sets of the remaining 362 PDB IDs are presented in Figure 4.

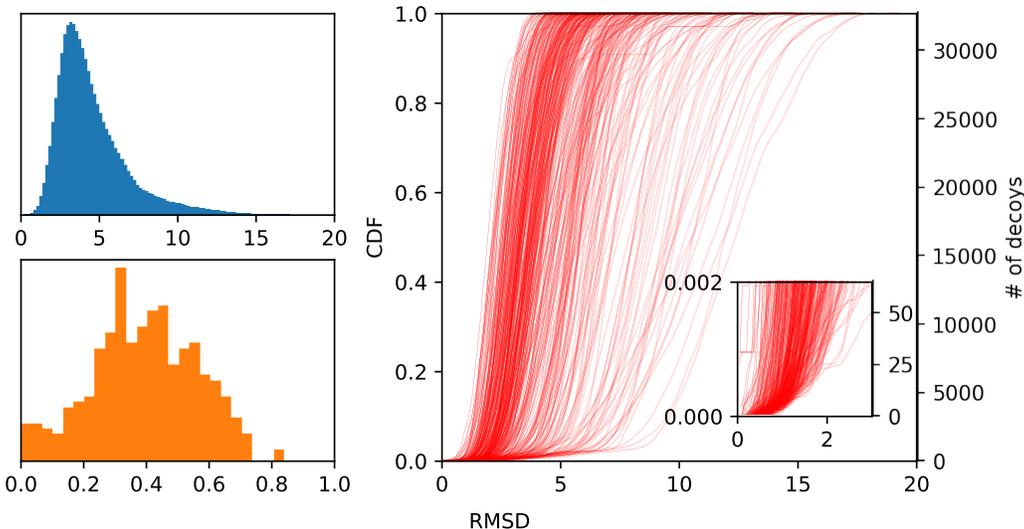


Figure 4: **Decoy set statistics.** *Top left:* histogram of RMSDs across all decoy sets (12 M total). *Bottom left:* histogram of minimum RMSDs among the decoy sets. All sets have a pose less than RMSD < 1 Å from the true pose. *Right:* cumulative density function of RMSDs in each decoy set. *Bottom right inset:* all decoy sets have at least 23 poses with RMSD < 2 Å.

F.2 Hyperparameters

Our method involves hyperparameters at several levels.

- The learned scalar fields have 5 channels.
- To parameterize the scalar field (Equation 2), we use spherical harmonics up to $\ell = 2$ and 5 Gaussian RBFs evenly spaced from 0 Å to 5 Å.
- All translational Fourier coefficients (Equation 4) are evaluated with a grid of frequencies corresponding to a sampling interval of 1 Å and a cubical domain with side length 40 Å. The integral over \mathbb{R}^3 in the training objective (Equation 9) is computed only over the cubical

FFT domain. During training and default inference, the cross-correlation is also computed with a sampling interval of 1 Å, but denser sampling intervals at inference-time (i.e., by zero-padding in the Fourier domain) are explored in Appendix G.

- Global spherical harmonic expansions (Equation 5) are computed up to $\ell = 10$, with 25 Gaussian RBFs evenly spaced from 0 Å to 20 Å. During training and default inference, the evaluation of rotational cross-correlations with FFTs (Equation 7) is always performed up to $\ell = 50$ and $\ell = 25$, respectively, by zero-padding in the Fourier domain, with other inference orders at inference-time explored in Appendix G.
- Local-to-global transformation matrices were precomputed for discretized positions along the $+z$ axis from 0 Å to 20 Å at 1 Å intervals.
- Data featurization and model hyperparameters are adapted from the default settings of Corso et al. (2023), giving a model size of 2.2 M parameters for both the ligand and protein model.
- By default, in the **RF** procedure, we evaluate $9^3 = 729$ translational grid points at inference time, filling a 8 Å cube at 1 Å intervals. In the **TF** procedure, we use an $m = 2$ grid over $SO(3)$ as implemented by Zhong et al. (2019) and Yershova et al. (2010), yielding 4608 grid points. Other resolutions are explored in Appendix G.

These hyperparameters were not extensively tuned, and further tradeoffs and improvements in performance and runtime could be explored by modifying them.

F.3 Runtime Measurements

All runtime measurements were performed on a machine with 64 Intel Xeon Gold 6130 CPUs and 8 Nvidia Tesla V100 GPUs. Gnina was run with default thread count settings. All of our processes were run on a single V100 GPU. For our method, we performed runtime analysis using CUDA events to remove the effects of asynchronous CUDA execution. Script loading, model loading, and algorithmic-level precomputations (which, if necessary, can be cached on disk) were excluded from the analysis. For Gnina, we attempted to remove similar overhead by timing single-pose scoring-only runs as representative of constant overhead costs. We report conformer docking runtimes in Table 3 using the PDBBind crystal structures; ESMFold runtimes are marginally shorter. Typical runtimes reported in Table 1 and Appendix D are obtained from timing runs with our method across the entire PDBBind test set.

F.4 Datasets

As noted previously, we use train, validation, and test splits from Stärk et al. (2022). However, due to RDKit parsing issues with Gnina-docked poses, the following 30 complexes are excluded (leaving 333 remaining) from all rigid conformer docking comparisons against Gnina, i.e., Tables 3 and Appendix G: 6HZB, 6E4C, 6PKA, 6E3P, 6OXT, 6OY0, 6HZA, 6E6W, 6OXX, 6HZZ, 6K05, 6NRH, 6OXW, 6RTN, 6D3Z, 6HLE, 6PY0, 6OXS, 6E3O, 6HZC, 6Q38, 6E7M, 6OIE, 6D3Y, 6D40, 6UHU, 6CJP, 6E3N, 6Q4Q, 6D3X. Scoring comparisons include all test complexes except 6A73, for which decoy poses could not be generated.

We download the 77 PDB IDs provided in Tosstorff et al. (2022) from the PDB to form the PDE10A dataset, keeping the A chain of each asymmetric unit and the Ligand of Interest (LOI) interacting with it. We then align all ligands to the crystal structure of 5SFS using the procedure described in Corso et al. (2023) for aligning ESMFold structures, except transforming the ligand rather than the protein. This constitutes the construction of a *cross-docking* dataset due to the use of the same pocket for all ligands. Due to RDKit parsing errors with the Gnina-docked poses, the following 7 PDB IDs are excluded from all comparisons: 5SFA, 5SED, 5SFO, 5SEV, 5SF9, 5SDX, 5SFC. The remaining 70 ligands are shown superimposed on the 5SFS pocket in Figure 5.

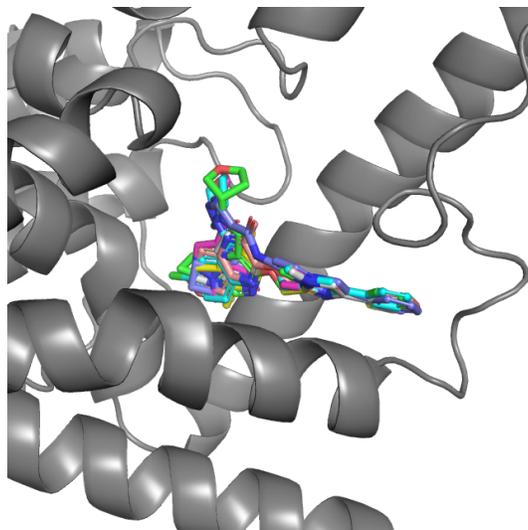


Figure 5: **PDE10A ligands** aligned on 5SFS

G Further Results

In Tables 4–7 below, we explore the impact of inference-time hyperparameters on the performance and runtime of our method on the rigid conformer docking task. We use the **ESF-N** model variant and experiment with the PDBBind crystal test set and PDE10A test set. For the **TF** procedure (Algorithm 1), we adjust (1) the number of grid points over $SO(3)$ with two possible resolutions, giving 576 and 4608 possible rotations, respectively; (2) the spatial interval at which the translational cross-correlation (Equation 3) is evaluated: either 1 Å or 0.5 Å (by zero-padding the scalar fields in the Fourier domain), giving a cubical grid of evaluated points with side length 8 Å and 9 or 17 points on each side. For the **RF** procedure (Algorithm 2), we adjust (1) the number of translational grid points with three possible resolutions, filling the 8 Å cube with 7, 9, or 13 points per side, respectively; (2) the resolution at which the rotational cross-correlation (Equation 6) is evaluated; this can be adjusted by zero-padding in the Fourier domain to include larger values of ℓ . The rows corresponding to results in the main Table 3 are bolded.

In all rows, the effective number of poses searched over via both degrees of freedom is computed. To provide an idea of the impact of discretization, we compute the median RMSD of the *closest grid point* to the ground-truth pose (decomposed into rotational and translational contributions). This serves as a hard lower bound for the median RMSD of the output docked pose. In the **TF** procedure, increasing the resolution is memory-intensive; thus, the **RF** procedure is more effective at leveraging FFT to conduct fine-grained search over the accelerated degree of freedom. The default reported performance is attained with a translational offset of 0.4 Å and a rotational offset of 0.16 Å. While performance improves with smaller grid offsets, the returns are rapidly diminishing.

The runtime of the method (averaged over 333 PDBBind complexes and 70 PDE10A complexes) is reported and color-coded according to Appendix D: protein preprocessing (green), ligand preprocessing (blue), and the pose optimization (red). The effect of the non-FFT grid resolution is also color-coded, i.e., in **TF** the explicit enumeration over $SO(3)$ grid points directly scales the ligand preprocessing, whereas in **RF** the enumeration over \mathbb{R}^3 scales the protein preprocessing. As the tables show, the preprocessing of these explicit grid points contributes to the majority of the non-amortizeable runtime. In general, the $SO(3)$ grid / ligand preprocessing in **TF** is less expensive, however, it cannot be amortized when moving from PDBBind to PDE10A (where the ligands are still distinct). On the other hand, the \mathbb{R}^3 grid / protein preprocessing time in **RF** is significantly reduced (very roughly on the order of 70-fold, as expected) in PDE10A compared to PDBBind.

Table 4: PDBBind TF

Trans. grid	SO(3) grid	Effective # poses	Grid offset			Med. RMSD	% <2Å	Runtime (ms)		
			Tr.	Rot.	All			Prot. prep.	Lig. prep.	Opt.
9	576	420k	0.52	0.84	0.98	1.53	63	65	931	100
9	4608	3.4M	0.50	0.42	0.67	1.10	72	72	7196	715
17	576	2.8M	0.25	0.80	0.84	1.50	64	70	928	123

Table 5: PDBBind RF

Trans. grid	SO(3) ℓ_{\max}	Effective # poses	Grid offset			Med. RMSD	% <2Å	Runtime (ms)		
			Tr.	Rot.	All			Prot. prep.	Lig. prep.	Opt.
7	10	3.2M	0.65	0.38	0.80	1.25	70	30k	85	158
7	25	45M	0.67	0.15	0.70	1.15	69	31k	87	225
7	50	353M	0.65	0.08	0.67	1.16	70	32k	85	704
9	10	6.8M	0.49	0.36	0.64	1.16	73	64k	85	333
9	25	97M	0.50	0.15	0.53	1.00	73	67k	87	476
9	50	751M	0.51	0.08	0.52	0.98	74	63k	84	1487
13	10	20M	0.33	0.37	0.51	1.05	74	198k	85	995
13	25	291M	0.33	0.15	0.37	0.90	72	200k	86	1430

Table 6: PDE10A TF

Trans. grid	SO(3) grid	Effective # poses	Grid offset			Med. RMSD	% <2Å	Runtime (ms)		
			Tr.	Rot.	All			Prot. prep.	Lig. prep.	Opt.
9	576	420k	0.51	0.88	1.00	1.85	56	22	761	89
9	4608	3.4M	0.50	0.48	0.69	1.11	64	21	6159	736
17	576	2.8M	0.26	0.89	0.93	2.05	50	20	756	106
17	4608	23M	0.26	0.44	0.51	1.00	73	20	6147	892

Table 7: PDE10A RF

Trans. grid	SO(3) ℓ_{\max}	Effective # poses	Grid offset			Med. RMSD	% <2Å	Runtime (ms)		
			Tr.	Rot.	All			Prot. prep.	Lig. prep.	Opt.
7	10	3.2M	0.72	0.38	0.83	1.60	54	476	44	161
7	25	45M	0.57	0.16	0.59	1.21	63	549	42	227
7	50	353M	0.65	0.08	0.65	1.30	64	635	59	718
9	10	6.8M	0.46	0.39	0.63	1.05	64	1014	42	327
9	25	97M	0.48	0.16	0.51	1.00	70	946	43	465
9	50	751M	0.49	0.09	0.50	0.99	64	943	42	1483
13	10	20M	0.34	0.41	0.55	1.17	64	2798	42	986
13	25	291M	0.33	0.16	0.36	0.96	69	2912	45	1469

